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RIGOROUS LOWER BOUND ON THE COMPRESSIBILITY OF A CLASSICAL SYSTEM

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A simple inequality is proved, from which a rigorous lower bound on the compressibility of a classical system with purely repulsive or hard core interaction follows immediately.

In this note we prove the following simple inequality: let p be a probability measure on the positive integers. Let $p_n = p_n/n$! be the probability of integer n. Suppose that the second moment of p exists and that for all $n \ge 0$:

$$(P_{n+2}/P_{n+1}) \ge (P_{n+1}/P_n) - A$$
, where $A > -1$. (a)

If some $P_{\mathcal{H}} = 0$, it is to be understood that $P_{\mathcal{H}} = 0$ for all $m \ge n$. Then:

$$\frac{\langle (n - \langle n \rangle)^2 \rangle}{\langle n \rangle} \ge \frac{1}{1 + A} \quad . \tag{1}$$

Equality in (1) is obtained:

if
$$A > 0$$
, for $P_n = \frac{N!}{(N-n)!} \frac{A^n}{(1+A)N}$ for $n \le N$,

(N arbitrary integer ≥ 1) and zero otherwise.

if
$$A = 0$$
, for $P_n = \alpha^n e^{-\alpha}$ with α real, $\alpha > 0$.

if -1 < A < 0, for $P_n =$ = $(-A)^n (1+A)^{\alpha} \alpha(\alpha+1) \dots (\alpha+n-1)$ with α real, $\alpha > 0$.

Proof: From Schwarz' inequality:

$$\{\langle n \rangle (1+A)\}^2 = \left\{ \sum_{0}^{\infty} \frac{1}{n!} P_n \left(\frac{P_{n+1}}{P_n} + nA \right) \right\}^2 \leq \sum_{0}^{\infty} \frac{1}{n!} P_n \left(\frac{P_{n+1}}{P_n} + nA \right)^2. (2)$$

expanding the right hand side and using (a), one gets (1). The cases of equality in (1) are obtained easily.

As an application, consider a system of clas-

sical particles, enclosed in a box of volume V, interacting through a 2 body potential φ ; the system is described in the grand canonical formalism. Let z be the fugacity, Z the grand partition function, T the absolute temperature, and β = = $(kT)^{-1}$. This defines a probability p on the integers by:

$$P_n = Z^{-1} z^n \int \mathrm{d}x_1 \dots \mathrm{d}x_n \exp[-\beta \sum_{i < j} \varphi(x_i - x_j)]. \tag{3}$$

For purely repulsive potentials and for hard core potentials, i.e. potentials such that (1): $\varphi(x-y) = +\infty$ for |x-y| < a, (2): for any finite sequence of n+1 points x_0, \ldots, x_n such that $|x_i - x_j| \ge a$ for all *i* and *j*,

$$\sum_{i=1}^{n} \varphi(x_0 - x_i) \ge -B,$$

where B is a real (positive) constant, one can prove easily [1,2] under the additional assumption

$$\int (1 - \exp[-\beta \varphi_+(x)]) \, \mathrm{d}x < \infty, \tag{b}$$

where $\varphi_+(x) = \max(\varphi(x), 0)$, that condition (a) is satisfied, with:

$$A \equiv zD \equiv z e^{\beta B} \int (1 - \exp[-\beta \varphi_{+}(x)]) dx$$
(4)

(For repulsive potentials, B = 0 and $\varphi_+ = \varphi_-$.) Therefore, for all positive z and V:

$$\chi \equiv \frac{1}{\rho} \frac{d\rho}{dP} \equiv \frac{\beta}{\rho} \frac{\langle (n - \langle n \rangle)^2 \rangle}{\langle n \rangle} \ge \frac{\beta}{\rho} \frac{1}{1 + \varepsilon D}$$

where $P = \ln Z/V$ is the pressure, $\rho = \langle n \rangle / V$ the

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density, and χ the compressibility. For repulsive or hard core potentials satisfying a weak tempering condition which is closely related to (b) and slightly stronger, the thermodynamic limit exists [3]. Inequality (5) still holds in the limit as a Lipschitz condition on the pressure. This proves directly in the grand canonical formalism the continuity of the pressure as a function of the density (proofs have already been given in the canonical formalism under various assumptions [1,2,4]) and gives a lower bound on the compressibility.

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